# Measurement of the angle $\alpha$ at BABAR

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**Abstract.** We present recent measurements of the CKM angle  $\alpha$  using data collected by the BABAR detector at the PEP-II asymmetric-energy  $e^+e^-$  collider at the SLAC National Acce- lerator Laboratory, operating at the  $\Upsilon(4S)$  resonance. We present constraints on  $\alpha$  from  $B \to \pi\pi$ ,  $B \to \rho\rho$  and  $B \to \rho\pi$  decays.

#### 1. Introduction

The measurements of the angles  $\alpha$ ,  $\beta$  and  $\gamma$  of the Unitarity Triangle (UT) at the B-factories are providing precision tests of the Standard Model (SM) description of CP violation. This description is provided by the Cabibbo-Kobayashi-Maskawa (CKM) quark-mixing [1, 2]. We summarize the experimental constraints on the  $\alpha$  UT angle obtained from B-meson decays to  $\pi\pi$ ,  $\rho\rho$  and  $\rho\pi$  with the BABAR experiment at the SLAC National Accelerator Laboratory. The BABAR detector and the PEP-II accelerator are described elsewhere [3].

## 2. Analysis Method

#### 2.1. General formula

The decay of a neutral B-meson into a pair of  $\pi$  or  $\rho$  mesons,  $B \to hh$   $(h = \pi, \rho)$ , occurs via two topologies: a tree-level process and a one-loop penguin diagram. The CP parameter  $\lambda_{hh}$ , defined by  $\lambda_{hh} = \frac{p}{q} \overline{A}$ , where q and p are the complex coefficient that link the mass and the flavor eigenstates in the B system, and  $A(\overline{A})$  is the  $B^0(\overline{B}^0)$  decay amplitude, can be expressed in terms of  $\alpha$  as

$$\lambda_{hh} = e^{2i\alpha} \frac{1 - (|V_{td}^* V_{tb}|/|V_{ud}^* V_{ub}|)P/Te^{-i\alpha}}{1 - (|V_{td}^* V_{tb}|/|V_{ud}^* V_{ub}|)P/Te^{i\alpha}},$$
(1)

where T and P are complex amplitudes dominated by tree and penguin topologies, respectively. The quantity experimentally measured is the time-dependent decay rate

$$f_{Q_{\text{tag}}} = \frac{e^{-|\Delta t|/\tau}}{4\tau} \left[ 1 - Q_{\text{tag}} C_{hh} \cos(\Delta m_d \Delta t) + Q_{\text{tag}} S_{hh} \sin(\Delta m_d \Delta t) \right] , \qquad (2)$$

where  $\tau$  is the neutral B lifetime and  $\Delta m_d$  is the  $B^0\overline{B}^0$  oscillation frequency.  $\Delta t$  is the proper time difference between decays of the B to hh ( $B_{\rm rec}$ ), and the second B in the event, denoted by  $B_{\rm tag}$ . The  $Q_{\rm tag}$  parameter is related to the flavor of the  $B_{tag}$ :  $Q_{\rm tag} = +1(-1)$  if the  $B_{\rm tag}$  is a  $B^0$  ( $\overline{B}^0$ ). The CP-violating asymmetries  $C_{hh}$  and  $S_{hh}$  are related to the  $\lambda_{hh}$  parameter by

$$S_{hh} = 2\mathcal{I}m(\lambda_{hh})/(1+|\lambda_{hh}|^2)$$
,  $C_{hh} = (1-|\lambda_{hh}|^2)/(1+|\lambda_{hh}|^2)$ . (3)

 $S_{hh}$  reflects the CP-violation induced by the interference between the mixing and decay processes;  $C_{hh}$  is the direct CP-violating asymmetry which comes from the interference between different decay topologies. In the absence of penguin contributions (P=0),  $C_{hh}$  vanishes and  $S_{hh}$  is simply related to the CKM angle  $\alpha$  by  $S_{hh} = \sin(2\alpha)$ .

In the more general case of the  $B^0(\overline{B}^0) \to \rho^{\pm} \pi^{\mp}$  decays, the time-dependent decay rate is given by

$$f_{Q_{\text{tag}}}^{\rho^{\pm}\pi^{\mp}} = (1 \pm \mathcal{A}_{\rho\pi}) \frac{e^{-|\Delta t|/\tau}}{4\tau} \left[ 1 - Q_{\text{tag}}(C_{\rho\pi} \pm \Delta C_{\rho\pi}) \cos(\Delta m_d \Delta t) + Q_{\text{tag}}(S_{\rho\pi} \pm \Delta S_{\rho\pi}) \sin(\Delta m_d \Delta t) \right] , \tag{4}$$

where, the  $\pm$  sign depends on whether the  $\rho$  meson is emitted by the W boson or comes from the spectator quark.  $\mathcal{A}_{\rho\pi}$  is the direct CP violation parameter measuring the asymmetry between the  $\rho^+\pi^-$  and  $\rho^-\pi^+$  final states, while  $\Delta S_{\rho\pi}$  and  $\Delta C_{\rho\pi}$ , which arise from the fact that two production modes of the  $\rho$  are possible, are dilution terms and have no CP content.

### 2.2. The isospin analysis

Using the strong isospin symmetry, the angle  $\alpha$  can be extracted up to discrete ambiguities from the CP-violating asymmetries defined above [4]. The decay amplitudes of the isospin-related final states obey the pentagonal relations

$$\sqrt{2}(A_{\rho\pi}^{+0} + A_{\rho\pi}^{0+}) = 2A_{\rho\pi}^{00} + A_{\rho\pi}^{+-} + A_{\rho\pi}^{-+} , \qquad \sqrt{2}(\overline{A}_{\rho\pi}^{+0} + \overline{A}_{\rho\pi}^{0+}) = 2\overline{A}_{\rho\pi}^{00} + \overline{A}_{\rho\pi}^{+-} + \overline{A}_{\rho\pi}^{-+} ; \qquad (5)$$

where  $A_{\rho\pi}^{ij}=A(B^0 \text{ or } B^+\to \rho^i\pi^j)$  and  $\overline{A}_{\rho\pi}^{ij}=A(\overline{B}^0 \text{ or } B^-\to \rho^i\pi^j)$ , i,j=+,-,0. With the use of these relations, 12 unknowns (6 complex amplitudes with one unphysical phase, and the CKM angle  $\alpha$ ) are to be determined while 13 observables are available:  $S_{\rho\pi}$ ,  $C_{\rho\pi}$ ,  $\Delta S_{\rho\pi}$ ,  $\Delta C_{\rho\pi}$ ,  $A_{\rho\pi}$ ; four average branching fractions  $\mathcal{B}(B\to \rho\pi)$ ; two time-dependent CP-violating asymmetries in the  $B^0\to \rho^0\pi^0$  decay  $(S_{\rho\pi}^{00},C_{\rho\pi}^{00})$  and two direct CP asymmetries in  $B^+\to \rho^+\pi^0$  and  $B^+\to \rho^0\pi^+$  decays.

In the case of  $B \to hh$   $(h = \pi, \rho)$ , Eq. 5 simplify to triangular relations

$$\sqrt{2}A_{hh}^{+0} = A_{hh}^{+-} + A_{hh}^{00} , \qquad \sqrt{2}\overline{A}_{hh}^{+0} = 2\overline{A}_{hh}^{+-} + \overline{A}_{hh}^{00} . \tag{6}$$

The information counting leads then to 6 unknowns and 7 observables: 3 branching fractions  $\mathcal{B}(B \to hh)$ ;  $C_{hh}$ ,  $S_{hh}$ ,  $C_{hh}^{00}$ ,  $S_{hh}^{00}$ . In the  $\pi\pi$  system  $S_{\pi\pi}^{00}$  is impossible to measure (as the  $\pi^0$  is reconstructed from two-photons decays, there is no way to measure the decay vertex), then one is left with 6 observables:  $\alpha$  can be extracted with an 8-fold ambiguity within  $[0, \pi]$  [5].

## 3. Experimental Results

# 3.1. $B \rightarrow \pi\pi$ and $B \rightarrow \rho\rho$

The various branching fractions and CP-asymmetries measured in  $B \to \pi\pi$  and  $B \to \rho\rho$  decays are summarized in Table 1. In the case of charged decays the charge asymmetry is defined as  $\mathcal{A}_{CP}(B \to hh) = -C_{hh}$ . The measurements are sufficiently well established to perform an isospin analysis.

The present measurement for the  $\pi^+\pi^-$  mode excludes the absence of CP violation  $(C_{\pi\pi}, S_{\pi\pi}) = (0,0)$  at a C.L. of 6.7 $\sigma$ . The relatively high branching fraction of the  $\pi^0\pi^0$  mode tends to separate the 8-fold ambiguities in the  $\alpha$  extraction, which only allows a weak constraint on  $\alpha$  to be set. With the current experimental measurements two of the eight ambiguities are nearly merged. The range [23°, 67°] in  $\alpha$  is excluded at the 90% C.L. [6]. The solution is in agreement with the global CKM fit [13, 14] which gives the range [71°, 109°] at 68% C.L.

Mode	$\mathcal{B}(10^{-6})$	C	S
$\pi^+\pi^-$	$5.5 \pm 0.4 \pm 0.3$ [7]	$-0.68 \pm 0.10 \pm 0.03$ [6]	$-0.25 \pm 0.08 \pm 0.02$ [6]
$\pi^0\pi^0$	$1.83 \pm 0.21 \pm 0.13$ [6]	$-0.43 \pm 0.26 \pm 0.05$ [6]	_
$\rho^+ \rho^-$	$25.5 \pm 2.1^{+3.6}_{-3.9}$ [9]	$0.01 \pm 0.15 \pm 0.06$ [9]	$-0.17 \pm 0.2^{+0.05}_{-0.06}$ [9]
$\rho^0 \rho^0$	$0.92 \pm 0.32 \pm 0.14$ [10]	$0.2 \pm 0.8 \pm 0.3$ [10]	$0.3 \pm 0.7 \pm 0.2$ [10]
Mode	$\mathcal{B}(10^{-6})$	$\mathcal{A}_{\mathbf{CP}}$	
$\pi^{\pm}\pi^{0}$	$5.02 \pm 0.46 \pm 0.29$ [8]	$0.03 \pm 0.08 \pm 0.01$ [8]	
$ ho^{\pm} ho^0$	$23.7 \pm 1.4 \pm 1.4$ [11]	$-0.054 \pm 0.055 \pm 0.010$ [11]	

**Table 1.** Summary of BABAR measurements of  $B \to \pi\pi$  and  $B \to \rho\rho$  decays. The measurements for the  $\rho\rho$  system corresponds to the longitudinal component of the decay rate. The errors quoted are statistical and systematic, respectively.

The analysis of  $B\to\rho\rho$  is potentially complicated due to the possible presence of three helicity states for the decay. The helicity zero state, which corresponds to longitudinal polarization of the decay, is CP-even but the helicity  $\pm 1$  states are not CP eigenstates. Fortunately this complication is avoided by the experimental finding that the dominant polarization is longitudinal,  $f_L(\rho^+\rho^-)=0.992\pm0.024^{+0.026}_{-0.013}$  [9],  $f_L(\rho^0\rho^0)=0.75^{+0.11}_{-0.14}\pm0.05$  [10] and  $f_L(\rho^+\rho^0)=0.950\pm0.015\pm0.006$  [11] ( $f_L\equiv\Gamma_L/\Gamma$ , where  $\Gamma$  is the total decay rate and  $\Gamma_L$  is the rate of the longitudinally-polarized mode). The  $B^0\to\rho^0\rho^0$  branching fraction is small compared with that of the  $B^+\to\rho^+\rho^0$  mode, which indicates that the penguin to three ratio (P/T, cf. Eq. 1) is small compared with that of the  $B\to\pi\pi$  system [4]. This has the effect of merging the different ambiguities in the extraction of  $\alpha$ . The latest  $B^0\to\rho^0\rho^0$  BABAR results present the first measurement of the time-dependent CP asymmetries  $C_L^{00}$  and  $S_L^{00}$ . The inclusion of these measurements has the effect of raising the 8-fold degeneracy on  $\alpha$ : the data only favors two solutions out of eight [10, 11]. These two effects allow to set a strong constraint on  $\alpha$ , where only two solutions are seen, corresponding to  $\alpha=(92.4^{+6.0}_{-6.5})^o$  at 68% C.L. [11] for the one in agreement with the global CKM fit [13, 14].

## 3.2. $B \rightarrow \rho \pi$

The  $B \to \rho \pi$  measurement reported here is a time-dependent amplitude analysis of  $B^0 \to (\rho \pi)^0$ . The interferences between the intersecting  $\rho$  resonance bands are modeled over the whole Dalitz Plot using the isobar model [15]. This allows determination of the strong phase differences from the interference pattern, which permits direct extraction of the angle  $\alpha$  with reduced ambiguities. The Dalitz amplitudes and time-dependence are contained in the 26 coefficients of the bilinear form-factor terms occurring in the time-dependent decay rate, which are determined from a likelihood fit. The values obtained for these coefficients are converted back into the quasi-two-body CP observables (c.f. Eq. 4), which are more intuitive in their interpretation. Table 2 reports the experimental findings on these observables [12].

Observable	Value	Observable	Value
$C_{\rho\pi}$ $\Delta C$	$0.15 \pm 0.09 \pm 0.05$ $0.39 \pm 0.09 \pm 0.09$	$S_{ ho\pi} \over \Lambda S$	$-0.03 \pm 0.11 \pm 0.04$ $-0.01 \pm 0.14 \pm 0.06$
$C^{00}_{ ho\pi}$	$-0.10 \pm 0.40 \pm 0.53$	$S_{ ho\pi}^{00}$	$0.01 \pm 0.14 \pm 0.00$ $0.04 \pm 0.44 \pm 0.18$
$A_{\rho\pi}^{\rho\pi}$	$-0.14 \pm 0.05 \pm 0.02$	p n	

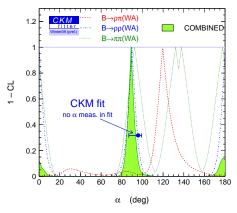
**Table 2.** Summary of BABAR measurements from the time-dependent amplitude analysis of  $B^0 \to (\rho \pi)^0$  decays. The errors quoted are statistical and systematic, respectively.

These measurements allow the determination of the limit  $\alpha = (87^{+45}_{-13})^o$  at 68% C.L., with almost no constraint at 95% C.L. This result is particularly interesting as there is an unique

solution in the  $[0, 180]^o$  range, which helps to break the ambiguities obtained from the  $\pi\pi$  and  $\rho\rho$  results. A hint of CP-violation is obtained at the level of  $3\sigma$ .

## 4. Summary

Several analyses have been conducted in BABAR to extract the angle  $\alpha$  of the UT. In the last few years the measurements of this angle have become increasingly precise. The measurements provided from the  $B \to \pi \pi/\rho \rho/\rho \pi$  modes give complementary constraints on  $\alpha$ . For the  $B \to \rho \rho$  system, the inclusion of the  $S_{\rho\pi}^{00}$  observable allows to favor two of the 8-fold ambiguities on  $\alpha$ , and the relatively large  $\mathcal{B}(B^+ \to \rho^+ \pi^0)$ , with respect to  $\mathcal{B}(B^0 \to \rho^0 \pi^0)$ , causes the ambiguities to degenerate in two peaks, improving the precision of the constraint. The measurements from the  $B^0 \to (\rho \pi)^0$  time-dependent amplitude analysis give a direct access to  $\alpha$ , disfavoring the ambiguities. The combined constraint averaging all the  $\pi \pi$ ,  $\rho \rho$  and  $\rho \pi$  measurements from BABAR and Belle gives  $\alpha = (89.0^{+4.4}_{-4.2})^o$  at 68% C.L. (see Fig. 1), which is in good agreement with the global CKM fit [13, 14].



**Figure 1.** Constraints on  $\alpha$ , provided by the CKMfitter group [13], expressed as one minus the confidence level as a function of angle. The constraints are constructed averaging the BABAR and Belle measurements for the  $\pi\pi$  (dotted green curve),  $\rho\rho$  (dash-dotted blue curve) and  $\rho\pi$  (dashed red curve) systems. The solid filled green curve represents the combined constraint using all the systems.

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# 6. References

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